Unitary version of the classical hamiltonians of the Fe_4 and Fe_8 SMMs

This text is part of sections 2 and 4 of Ref. [1].

The main terms in the hamiltonian of the Fe_4 and Fe_8 single-molecule magnets (SMMs) are of the same type. Their hamiltonian in the presence of an external constant magnetic field with arbitrary direction is [2, 3]

$$\mathbf{H}_{S} = \sum_{j=1}^{N} \mathbf{H}_{j}^{(S)}, \qquad j = 1, 2, \cdots, N,$$
 (1a)

where

$$\mathbf{H}_{j}^{(S)} = -D_{S}(S_{j}^{z})^{2} + E_{S}[(S_{j}^{x})^{2} - (S_{j}^{y})^{2}] - h_{x}S_{j}^{x} - h_{y}S_{j}^{y} - h_{z}S_{j}^{z}.$$
 (1b)

The hamiltonian (1a) describes a Fe₄ or Fe₈ SMM with N sites, in which h_x , h_y and h_z are the components of the constant external magnetic field. For the Fe₄ SMM, (S_j^x, S_j^y, S_j^z) corresponds to spin-5 operators at the *j*th-site, whereas for the Fe₈ SMM they are the spin-10 operators.

The classical versions of these two SMMs are obtained by replacing, in (1b), the spin-S operators by the classical functions

$$S_x^{class} = \sqrt{S(S+1)} \sin\theta \cos\phi$$
 (2a)

$$S_y^{class} = \sqrt{S(S+1)} \sin \theta \sin \phi$$
 (2b)

$$S_x^{class} = \sqrt{S(S+1)} \cos\theta , \qquad (2c)$$

with $\theta \in [0, \pi]$ and $\phi \in [0, 2\pi]$.

An interesting comparison of some thermodynamical functions of the Fe₄ and Fe₈ SMMs can be made when we study the thermodynamics of the unitary (normalized) version of the hamiltonian (1b) for the spins S = 5 and S = 10. For each value os spin in this hamiltonian, the unitary version of $\mathbf{H}_{j}^{(S)}$ is obtained by replacing the spin matrices by the unitary spin operator \vec{s} , defined as

$$\vec{s}_j \equiv \frac{\vec{S}_j}{\sqrt{S(S+1)}}, \quad S \in \{5, 10\},$$
(3)

where \vec{S}_j is the spin operator at the *j*th- site, with norm $||\vec{S}_j|| = \sqrt{S(S+1)}$, for $S \in \{5, 10\}$.

The unitary version of the hamiltonian (1b) is

$$\mathcal{H}_{j}^{(S)} = -D_{S}(s_{j}^{z})^{2} + E_{S}[(s_{j}^{x})^{2} - (s_{j}^{y})^{2}] - h_{x}s_{j}^{x} - h_{y}s_{j}^{y} - h_{z}s_{j}^{z}.$$
(4)

We denote by $\omega_S(D_S, E_S, h_x, h_y, h_z; \beta)$ the Helmholtz free energy (HFE) derived from (4) for S = 5 and S = 10.

It is simple to show that the HFEs derived from the hamiltonians (1b) (we call it $\mathcal{W}_S(D_S, E_S, h_x, h_y, h_z; \beta)$) and (4) are related by

$$\mathcal{W}_S(D_S, E_S, h_x, h_y, h_z; \beta) = S(S+1) \,\omega_S(D_S, E_S, \tilde{h}_x, \tilde{h}_y, \tilde{h}_z; \beta).$$
(5)

where

$$\tilde{h}_i \equiv \frac{h_i}{\sqrt{S(S+1)}}, \qquad \tilde{\beta} \equiv S(S+1)\beta$$
(6)

and $i \in \{x, y, z\}$.

References

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